

**Скрытая масса
как следствие теоремы вириала Пуанкаре**

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$$m_i \frac{d^2 \vec{r}_i}{dt^2} = \vec{F}_i \quad (i=1, 2, \dots, n), \dots (1)$$

$$m_i \left(\vec{r}_i \cdot \frac{d^2 \vec{r}_i}{dt^2} \right) = \left(\vec{r}_i \cdot \vec{F}_i \right), \dots (2)$$

$$\sum_{i=1}^n m_i \left(\vec{r}_i \cdot \frac{d^2 \vec{r}_i}{dt^2} \right) = \sum_{i=1}^n \left(\vec{r}_i \cdot \vec{F}_i \right), \dots (3)$$

$$\left(\vec{r}_i \cdot \frac{d^2 \vec{r}_i}{dt^2} \right) = \frac{d}{dt} \left(\vec{r}_i \cdot \frac{d \vec{r}_i}{dt} \right) - \left(\frac{d \vec{r}_i}{dt} \right)^2, \dots (4)$$

$$\left(\vec{r}_i \cdot \frac{d^2 \vec{r}_i}{dt^2} \right) = \frac{1}{2} \frac{d}{dt} \left(\vec{r}_i^2 \right), \dots (5)$$

$$\left(\vec{r}_i \cdot \frac{d^2 \vec{r}_i}{dt^2} \right) = \frac{1}{2} \frac{d^2 \vec{r}_i}{dt^2} - \left(\frac{d \vec{r}_i}{dt} \right)^2, \dots (6)$$

$$\frac{1}{2} \sum_{i=1}^n m_i \frac{d^2 \vec{r}_i}{dt^2} - \sum_{i=1}^n m_i \left(\frac{d \vec{r}_i}{dt} \right)^2 = \left(\vec{r}_i \cdot \vec{F}_i \right) \dots (7)$$

$$J = \sum_{i=1}^n m_i \vec{r}_i^2, \dots (8)$$

$$T = \frac{1}{2} \sum_{i=1}^n m_i \left(\frac{d \vec{r}_i}{dt} \right)^2, \dots (9)$$

$$V = \sum_{i=1}^n \left(\vec{r}_i \cdot \vec{F}_i \right) \dots (10)$$

$$\frac{1}{2} \frac{d^2 J}{dt^2} = V + 2T \dots (11)$$

$$1). \frac{d^2 J}{dt^2} = 0 \dots (12)$$

$$\bar{T} = -\frac{1}{2} \bar{V} \dots (13)$$

(теорема Виршала, Крауфорд, 1870)

$$\vec{F}_i = -\frac{\partial \Omega}{\partial \vec{r}_i}, \quad \Omega = -G \sum_{i=2}^n \sum_{j=1}^{i-1} \frac{m_i m_j}{r_{ij}}, \quad r_{ij} = |\vec{r}_j - \vec{r}_i| \dots (14)$$

$$\Omega(l \vec{r}_i, l \vec{r}_j) = l^{-1} \Omega(\vec{r}_i, \vec{r}_j), \quad \sum_{i=1}^n \left(\vec{r}_i \cdot \frac{\partial \Omega}{\partial \vec{r}_i} \right) = -\Omega \dots (15)$$

$$V = -\sum_{i=1}^n \left(\vec{r}_i \cdot \frac{\partial \Omega}{\partial \vec{r}_i} \right) = +\Omega \dots (16) \quad 2\bar{T} = -\bar{\Omega} \text{ (Пугачев, 1913)}$$

\dots (17)

$$2). \frac{d^2 J}{dt^2} \neq 0 \dots (18),$$

$$\frac{1}{2} \frac{d^2 J}{dt^2} = 2\bar{T} + \bar{\Omega} \dots (19)$$

(Эггерстон, 1916)

$$(20) \dots R_{ij} - \frac{1}{2} g_{ij} R - \Lambda g_{ij} = \kappa T_{ij}, \quad \kappa > 0, \Lambda > 0, (i, j = 0, 1, 2, 3)$$

$$(21) \dots R_{ij} = \frac{\partial \Gamma_{ij}^l}{\partial x^l} - \frac{\partial \Gamma_{lj}^i}{\partial x^l} + \Gamma_{ij}^l \Gamma_{lm}^m - \Gamma_{il}^m \Gamma_{jm}^l, \quad R = g^{ij} R_{ij}$$

$$(22) \dots \Gamma_{ij}^l = \Gamma_{m,ij} g^{lm}, \quad \Gamma_{m,ij} = \frac{1}{2} \left(\frac{\partial g_{mi}}{\partial x^j} + \frac{\partial g_{mj}}{\partial x^i} - \frac{\partial g_{ij}}{\partial x^m} \right)$$

$$(23) \dots g_{ij} = g_{ji}, \quad g^{il} g_{lk} = \delta_k^i, \quad \delta_k^i = \begin{cases} 1 & (i=k) \\ 0 & (i \neq k) \end{cases}$$

$$(24) \quad ds^2 = g_{ik} dx^i dx^k, \quad x^0 \equiv ct, \quad x^1 \equiv r, \quad x^2 \equiv \theta, \quad x^3 \equiv \varphi$$

(А. Эйнштейн, 1916)

$$(25) \dots \begin{cases} g_{00} = 1 - \frac{2GM}{c^2 r} - \frac{1}{3} \Lambda r^2, & g_{ik} = 0 \quad (i \neq k) \\ g_{11} = \left(1 - \frac{2GM}{c^2 r} - \frac{1}{3} \Lambda r^2 \right)^{-1}, & g_{22} = -r^2, \quad g_{33} = -r^2 \sin^2 \theta \end{cases}$$

$$(26) \dots g_{00} = 1 - \frac{2\psi}{c^2}, \quad \psi = \frac{GM}{r} + \frac{1}{6} \Lambda c^2 r^2$$

$$(27) \dots \bar{W} = -\bar{\Omega} + \bar{B}$$

$$(28) \dots \bar{B} = \frac{1}{6} \Lambda c^2 \sum_{i=1}^n \sum_{j=1}^n m_i z_{ij}^2$$

$$(29) \dots \bar{\Omega} = -G \sum_{i=1}^n \frac{m_i^2}{R} \approx -G \frac{(nm)^2}{2R} = -\frac{GM^2}{2R}$$

$$(30) \dots M = \sum_{i=1}^n m_i = n \cdot m$$

$$(31) \dots \bar{B} = \frac{1}{6} \Lambda c^2 M R^2$$

$$(32) \dots \langle v^2 \rangle = \frac{1}{M} \sum_{i=1}^n m_i \left(\frac{d\vec{z}_i}{dt} \right)^2$$

$$(33) \dots \frac{1}{2} \frac{d^2 J}{dt^2} = (\bar{\Omega} + 2\bar{B}) + 2\bar{T}$$

$$(34) \dots \frac{d^2 J}{dt^2} = 0, \quad 2\bar{T} = -\bar{\Omega} - 2\bar{B}$$

$$(35) \dots M \langle v^2 \rangle = + \frac{GM^2}{2R} - \frac{1}{6} \Lambda c^2 M R^2$$

$$(36) \quad \langle v^2 \rangle = \frac{GM}{2R} - \frac{1}{6} \Lambda c^2 R^2$$

$$(37) \quad M = \frac{\Lambda c^2}{3G} R^3 + \frac{2}{G} R \langle v^2 \rangle$$

$$(38) \quad \rho_v = \frac{\Lambda c^2}{8\pi G}, \quad M = \frac{2R \langle v^2 \rangle}{G} + \frac{8}{3} \pi \rho_v R^3$$

$$(39) \quad M = M_p + \Delta M, \quad M_p = \frac{2R \langle v^2 \rangle}{G}, \quad \Delta M = \frac{8}{3} \pi \rho_v R^3$$

$$(40) \quad \rho_G = -2\rho_v, \quad \Delta M = -\frac{4}{3} \pi R^3 \rho_G = M_v < 0.$$